



Some plans for Dynamics in HIRLAM-B

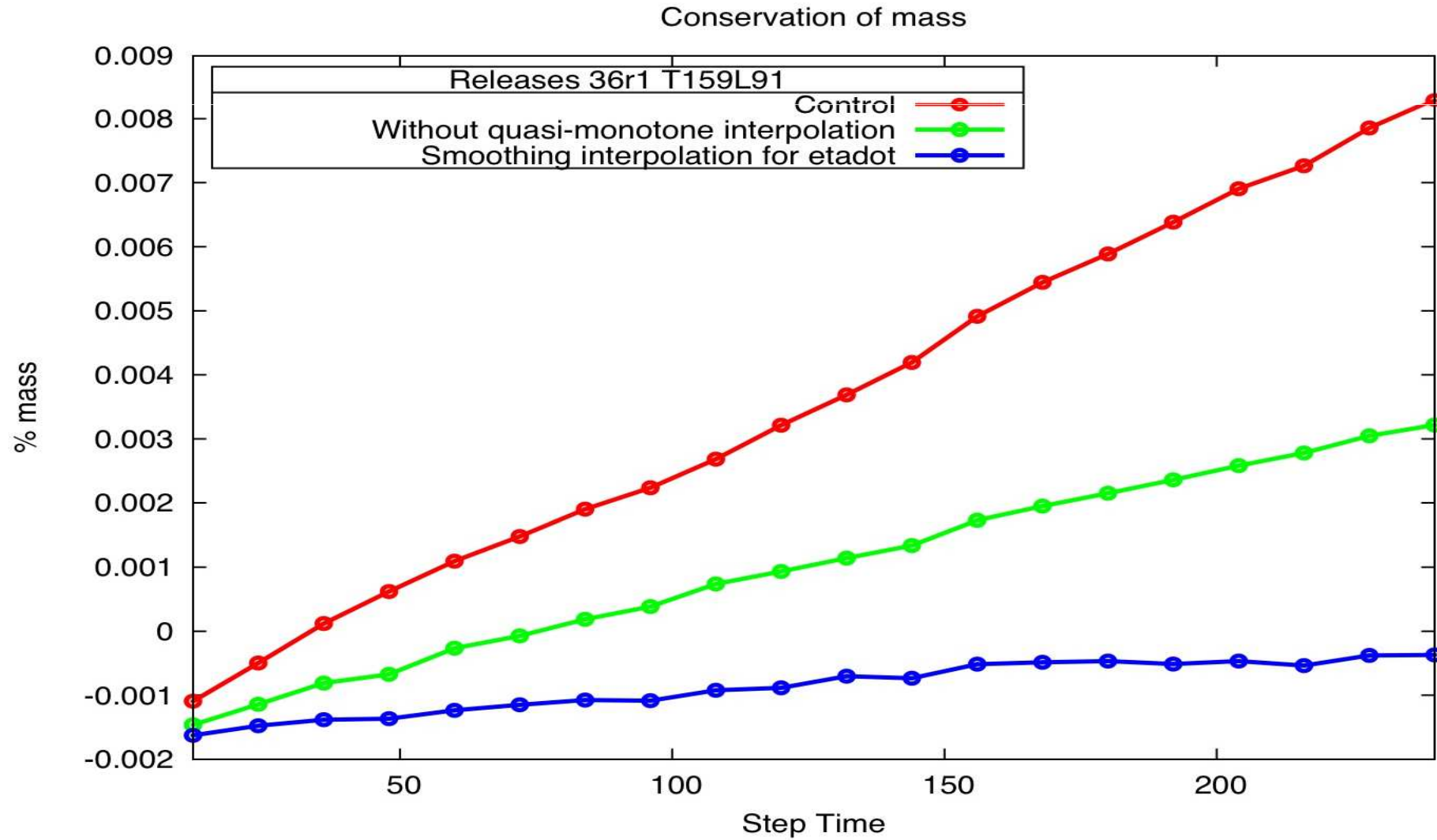
M. Hortal, PM for dynamics in
HIRLAM-A

Overview



- Improvement of mass conservation in semi-Lagrangian
- Physics-dynamics interface
 - Second-order accurate treatment of slow processes
 - Physics and dynamics on different grids
- Height-based vertical coordinate with finite elements
- Semi-analytical time stepping scheme
- Laplace transform time stepping scheme

Conservation in semi-Lagrangian



Physics-dynamics interface



- Second-order accurate treatment of the physical tendencies
 - Physical tendencies (at least the slow processes) treated as the non-linear terms (SETTLS)
- Physics and dynamics run on different grids
 - Bi-linear interpolation to go from the coarser to the finer grid
 - Area averaging to go from the finer to the coarser grid

Height based coordinate and covariant variables



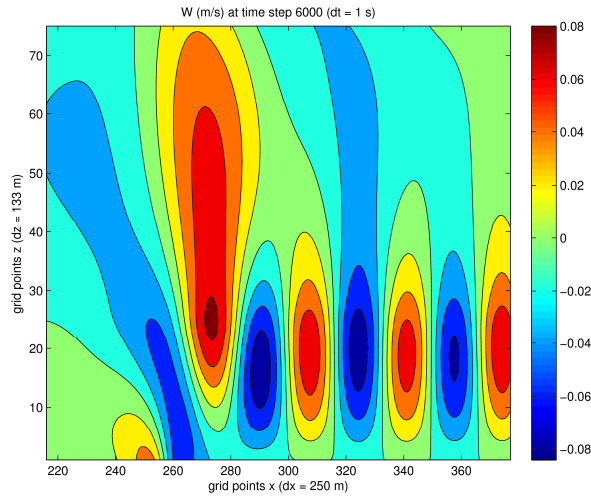
- Both mass and height based coordinates has been shown to be suitable for NH modeling
- Why to introduce height based coordinate in HARMONY?
 - Is a time independent coordinate system
 - Easier mathematical treatment of covariant variables
 - Covariant variables makes the expression of divergence operator simpler, without the non-linear term which is unstable in presence of orography
- Covariant vertical velocity variable W is the prognostic variable used in the Helmholtz structure equation of the semi-implicit solver
 - Robust boundary condition implementation: $W=0$ at boundaries is included in the semi-implicit solver.

2D model for testing

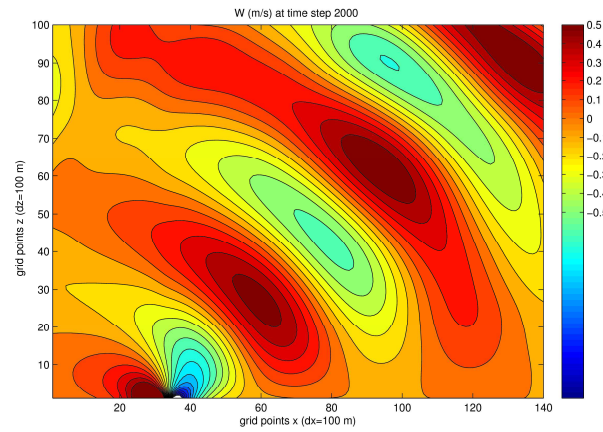


- SHB numerical linear stability analysis shows:
 - amplification factors under 1.01 in most cases: stable or marginally unstable.
 - stability for temperature T in the range $0.5 T^* < T < 1.5 T^*$.
 - In 2 time-levels stability with $T^* > T$
- Vertical discretization:
 - VFE has been implemented, is more accurate and slightly more stable than VFD.
- **Advection:**
 - Covariant Eulerian and covariant semi-Lagrangian schemes has been applied successfully.

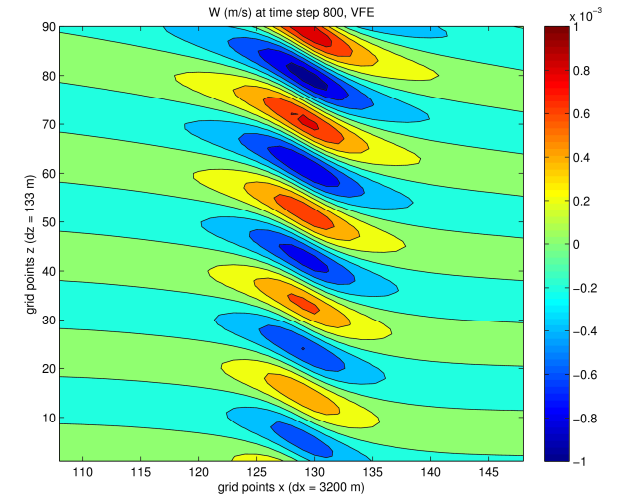
Tests done



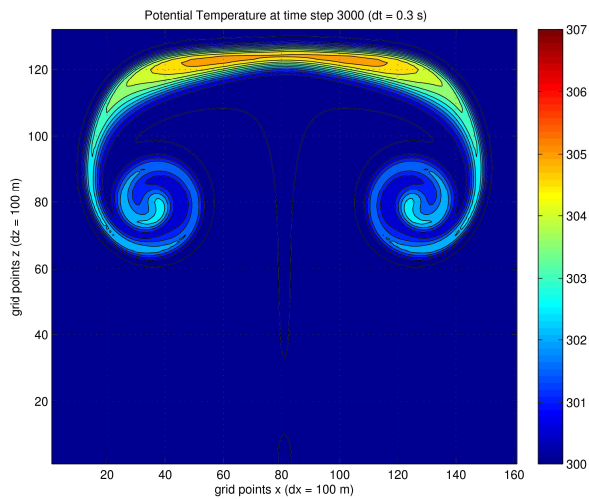
2 layer atmosphere with different stability parameters



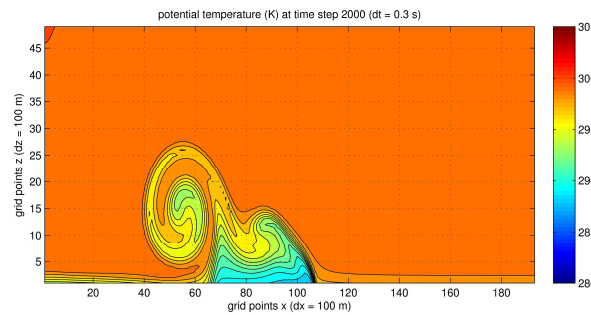
Quasi-linear non-hydrostatic flow



Hydrostatic flow

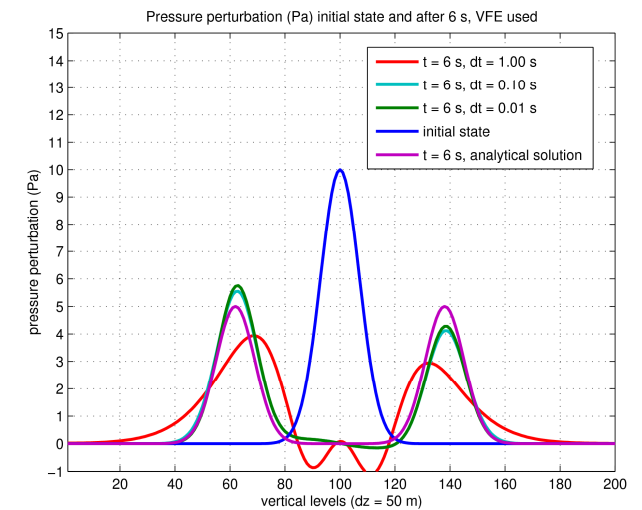


Warm bubble



Cold bubble

32 EWGLAM meeting, Exeter



Acoustic wave

Semi-analytical time stepping scheme



- Once the r.h.s. of the equations has been split between linear and non-linear terms
- Taking as constant (during the time-step) the non-linear terms
 - In spectral space and projected on the eigenvectors of the vertical linear operator
- The equations are ordinary linear differential equation which have analytical solution
- No need then to apply an implicit treatment !!!!

Semi-analytical time stepping scheme (cont)



- Example: Hydrostatic IFS
 - Linearized equations ignoring advection and Coriolis

$$\frac{\partial D}{\partial t} = -\gamma \nabla^2 T - R_d T_r \nabla^2 \ln p_s + NLD$$

$$\frac{\partial T}{\partial t} = -\tau D + NLT$$

$$\frac{\partial \ln p_s}{\partial t} = -\nu D + NLP$$

Semi-analytical time stepping scheme (cont)



Semi-implicit

$$D^+ = D^0 - \frac{\Delta t}{2} \gamma \nabla^2 T^+ - \frac{\Delta t}{2} R_d T_r \nabla^2 \ln p_s^+ - \frac{\Delta t}{2} \gamma \nabla^2 T^0 - \frac{\Delta t}{2} R_d T_r \nabla^2 \ln p_s^0 + NLD \cdot \Delta t$$

$$T^+ = T^0 - \frac{\Delta t}{2} \tau D^+ - \frac{\Delta t}{2} \tau D^0 + NLT \cdot \Delta t$$

$$\ln p_s^+ = \ln p_s^0 - \frac{\Delta t}{2} \nu D^+ - \frac{\Delta t}{2} \nu D^0 + NLP \cdot \Delta t$$

$$\left(1 - \frac{\Delta t^2}{4} \nabla^2 (\gamma \tau + R_d T_r \nu) \right) D^+ = \left(1 + \frac{\Delta t^2}{4} \nabla^2 (\gamma \tau + R_d T_r \nu) \right) D^0 - \Delta t \nabla^2 \left(\gamma T^0 + \frac{\Delta t}{2} \gamma NLT + \Delta t R_d T_r \ln p_s^0 - \frac{\Delta t}{2} R_d T_r NLP \right) + NLD \cdot \Delta t$$

Semi-analytical time stepping scheme (cont)



Semi-analytical

$$\begin{aligned}
 \frac{\partial^2 D}{\partial t^2} &= -\nabla^2 \gamma(-\tau D + NLT) - R_d T_r \nabla^2 (-\nu D + NLP) + \frac{\partial(NLD)}{\partial t} = \\
 &= \nabla^2 (\gamma\tau + R_d T_r \nu) D - \nabla^2 \gamma NLT - R_d T_r \nabla^2 NLP + \frac{\partial(NLD)}{\partial t} = \\
 &= \nabla^2 (\gamma\tau + R_d T_r \nu) D + NL
 \end{aligned}$$

Whose general solutions are

$$D(t) = Ae^{i\omega(t-t_0)} + Be^{-i\omega(t-t_0)} + \frac{NL}{\omega^2}$$

where $-\omega^2$ are the eigenvalues of $\nabla^2 (\gamma\tau + R_d T_r \nu)$

Semi-analytical time stepping scheme (cont)

At $t = t_0$

$$D(t_0) = A + B + \frac{NL}{\omega^2}$$

$$i\omega A - i\omega B = -\nabla^2 \gamma T(t_0) - R_d T_r \nabla^2 \ln p_s(t_0) + NLD$$

$$D(\Delta t) = D(0)\cos(\omega\Delta t) - \Delta t \left[\nabla^2 \gamma T(0) + R_d T_r \nabla^2 \ln p_s(0) - NLD \right] \frac{\sin(\omega\Delta t)}{\omega\Delta t} - \left(\nabla^2 \gamma NLT + R_d T_r \nabla^2 NLP \right) \frac{1 - \cos(\omega\Delta t)}{\omega^2}$$

Laplace transform time stepping scheme (from P. Lynch)



The Laplace transform of the general equation

$$\frac{d\mathbf{X}}{dt} + \mathbf{L} \cdot \mathbf{X} + N(\mathbf{X}) = 0$$

Gives, after rearranging

$$\hat{\mathbf{X}} = (s\mathbf{I} + \mathbf{L})^{-1} [\mathbf{X}^0 - N^0 / s]$$

Applying the inverse operator at $t = \Delta t$

gives the filtered state

$$\mathbf{X}(\Delta t) = \mathbf{L}^* \left\{ \hat{\mathbf{X}} \right\}_{t=\Delta t}$$